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ON SIMILARITY SOLUTIONS FOR THAWING PROCESSES*

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Abstract

We review some recent results for a mathematical model for thawing in a saturated semi-infinite porous medium when change of phase induces a density jump and a heat flux condition of the type $-q_0 t^{-\frac{1}{2}}$ is imposed on the fixed face x=0. Different cases depending on physical parameters are analysed and the explicit solution of the similarity type is obtained when a given condition for the thermal coefficient q_0 is verified.

Key words. Stefan problem, free boundary problems, phase change process, similarity solution, density jump, thawing processes, freezing, solidification.

Resumen. Se realiza una revisión de recientes resultados sobre un modelo matemático de descongelación en un medio poroso semi-infinito saturado cuando un cambio de fase induce un salto de densidad y cuando se impone una condición de flujo de calor del tipo $-q_0 t^{-\frac{1}{2}}$ en el borde fijo x=0. Se analizan diferentes casos que dependen de diversos parámetros físicos y se obtiene la solución explícita de tipo similaridad cuando una cierta condición sobre el coeficiente térmico q_0 es satisfecha.

Palabras Claves. Problema de Stefan, problemas de frontera libre, procesos de cambio de fase, solución de similaridad, salto de densidad, procesos de descongelación, congelación, solidificación.
AMS subject clasification. 35R35, 80A22, 35C05.

1 Introduction

Phase-change problems appear frequently in industrial processes and other problems of technological interest [8]. A large bibliography on the subject was given in [11]. In this paper, we consider the problem of thawing of a partially frozen porous media, saturated with an incompressible liquid, with the aim of constructing similarity solutions.

We have in mind the following physical assumptions (see [2], [4], [5]):

1. A sharp interface between the frozen part and the unfrozen part of the domain exists (sharp, in the macroscopic sense).

2. The frozen part is at rest with respect to the porous skeleton, which will be considered to be indeformable.

3. Due to density jump between the liquid and solid phase, thawing can induce either desaturation or water movement in the unfrozen region. We will consider the latter situation assuming that liquid is continuously supplied to keep the medium saturated.

Although thawing has received less attention than freezing, our investigation is in the same spirit as [3], and [9], with the simplification due to the absence of ice lenses and frozen fringes.

We will study a one-dimensional model of the problem, using the following notation:

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$\varepsilon > 0$	:	porosity,
$\rho > 0$	:	density; ρ_w and ρ_I: density of water and ice $\left(g/cm^3\right)$
$c > 0$	:	specific heat at constant density $\left(\frac{cal}{g^oC}\right)$
$k > 0$	:	conductivity $\left(\frac{cal}{s\,cm\,^oC}\right)$
u	:	temperature of unfrozen zone (^oC)
v	:	temperature of frozen zone (^oC)
$u = v = 0$	:	being the melting point at atmospheric pressure
$\lambda > 0$	:	latent heat at $u = 0$ (cal/g)
γ	:	coefficient in the Clausius-Clapeyron law ($s^2 cm^oC/g$)
$\mu > 0$	:	viscosity of liquid $\left(g/cm^3\right)$

and subscripts F, U, I and W refer to the frozen medium, unfrozen medium, pure ice and pure water, respectively, while S refers to the porous skeleton.

The unknowns of the problem are a function $x = s(t)$, representing the free boundary separating $Q_1 = \{(x,t) : 0 < x < s(t),\, t > 0\}$ and $Q_2 = \{(x,t) : s(t) < x,\, t > 0\}$, and the two functions $u(x,t)$ and $v(x,t)$ defined in Q_1 and Q_2, respectively. Besides standard requirements, $s(t)$, $u(x,t)$ and $v(x,t)$ fulfil the following conditions (we refer to [4] for a detailed explanation of the model):

$$u_t = a_1 u_{xx} - b\rho\dot{s}(t)\, u_x\,, \qquad \text{in } Q_1 \tag{1}$$
$$v_t = a_2 v_{xx}\,, \qquad \text{in } Q_2 \tag{2}$$
$$u(s(t),t) = v(s(t),t) = d\rho s(t)\,\dot{s}(t)\,, \qquad t > 0 \tag{3}$$
$$k_F v_x(s(t),t) - k_U u_x(s(t),t) = \alpha\dot{s}(t) + \beta\rho s(t)\,\dot{s}^2(t)\,, \qquad t > 0 \tag{4}$$
$$v(x,0) = v(+\infty,t) = -A < 0\,, \qquad x,t > 0 \tag{5}$$
$$s(0) = 0 \tag{6}$$
$$k_U u_x(0,t) = -\frac{q_0}{\sqrt{t}}, \qquad t > 0 \tag{7}$$

with

$$a_1 = \alpha_1^2 = \frac{k_U}{\rho_U c_U}, \qquad a_2 = \alpha_2^2 = \frac{k_F}{\rho_F c_F}, \qquad b = \frac{\varepsilon\rho_W c_W}{\rho_U c_U},$$
$$d = \frac{\varepsilon\gamma\mu}{K}, \qquad \rho = \frac{\rho_W - \rho_I}{\rho_W}, \qquad \alpha = \varepsilon\rho_I\lambda,$$
$$\beta = \frac{\varepsilon^2\rho_I\,(c_W - c_I)\,\gamma\mu}{K} = \varepsilon d\rho_I\,(c_W - c_I)\,.$$

Problem I consists of equations (1)-(7), while by problem II we mean the system and (1)-(6) and (8) respectively, where

$$u(0,t) = B > 0, \qquad t > 0\,. \tag{8}$$

Problem II was previously studied in [5].

We will look for similarity solutions of Problem I in different cases according to the value of parameters ρ, β y d, following the methods introduced in [5],[10].

First of all, we note that the function $u(x,t) = \Phi(\eta)$, with $\eta = \frac{x}{2\alpha_1\sqrt{t}}$, is a solution of (1) if and only if Φ satisfies the following equation

$$\frac{1}{2}\Phi''(\eta) + \left(\eta - \frac{b\rho}{\alpha_1}\sqrt{t}\dot{s}(t)\right)\Phi'(\eta) = 0$$

and similarly, the function $v(x,t) = \Psi(\eta)$ is solution of (2) if and only if Ψ satisfies the equation

$$\frac{1}{2}\Psi''(\eta) + \eta\Psi'(\eta) = 0.$$

Therefore, we obtain the following result

Theorem 1 *The free boundary problem I has the similarity solutions*

$$\begin{aligned} s(t) &= 2\xi\alpha_1\sqrt{t} \\ u(x,t) &= m\xi^2 + \frac{2q_0\alpha_1}{K_U} g(p,\xi) - \frac{2q_0\alpha_1}{K_U} \int_0^{\frac{x}{2\alpha_1\sqrt{t}}} \exp(pyr - r^2)\, dr \\ v(x,t) &= \frac{m\xi^2 + A \operatorname{erf}(\gamma_0\xi)}{\operatorname{erfc}(\gamma_0\xi)} - \frac{m\xi^2}{\operatorname{erfc}(\gamma_0\xi)} \operatorname{erf}\left(\frac{x}{2\alpha_1\sqrt{t}}\right) \end{aligned} \tag{9}$$

if and only if the coefficient ξ satisfies the equation

$$q_0 \exp\left((p-1)y^2\right) - K_2 F(m,y) = \delta y + \nu y^3, \qquad y > 0 \tag{10}$$

where

$$F(m,y) = (A + my^2)\frac{\exp(-\gamma_0^2 y^2)}{\operatorname{erfc}(\gamma_0 y)}, \qquad g(p,\xi) = \int_0^y \exp(pyr - r^2)\, dr \tag{11}$$

$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x \exp(-r^2)\, dr, \qquad \operatorname{erfc}(x) = 1 - \operatorname{erf}(x),$$

and the constants $K_2, m, \delta, \nu, \gamma_0$ are defined as follows

$$\begin{aligned} K_2 &= \frac{K_F}{\alpha_2\sqrt{\pi}} \qquad m = 2d\rho\alpha_1^2 \qquad \gamma_0 = \frac{\alpha_1}{\alpha_2} > 0 \\ \delta &= \alpha\alpha_1 > 0 \qquad \nu = 2\beta\rho\alpha_1^3 \qquad p = 2b\rho\,. \end{aligned} \tag{12}$$

In order to analize (10) we need some preliminary results.

Lemma 2 *(a) If $m > 0$, then F grows from A to $+\infty$, when y grows from 0 to $+\infty$. If $m < 0$, then F has a unique positive maximum, from which it decreases to $-\infty$. In both cases, $F(m,y) \sim \sqrt{\pi}\gamma_0 m y^3$ when $y \to +\infty$.*

(b) For all $p > 0$, we have

(i) $g(p,y) \geq \frac{1}{py}\left(\exp\left((p-1)\,y^2\right)\right) - \exp\left(-y^2\right), \qquad y > 0$

(ii) $g_y(p,y) \geq \exp\left((p-1)\,y^2\right) + \frac{p}{2}\left(1 - \exp\left(-y^2\right)\right) > 0, \qquad y > 0$

(iii) $g(p,0) = 0, \qquad g_y(p,0) = 1, \qquad g(p,+\infty) = +\infty.$

(iv) $\lim_{y\to+\infty} \frac{g(p,y)}{y^2} = 0$ *if* $p \leq 0$ *and* $\lim_{y\to+\infty} \frac{g(p,y)}{y^2} = +\infty$ *if* $p > 0$.

(v) $\lim_{y\to+\infty} \frac{y\, g(p,y)}{\exp((p-1)y^2)} = \begin{cases} \frac{1}{p-2} & \text{if } p > 2 \\ +\infty & \text{if } p \leq 2 \end{cases}$

(c) If $m > 0$ and $\nu \geq -m\sqrt{\pi}\gamma_0 K_2$, then the function

$$H(y) = F(m,y) + \frac{\nu}{K_2} y^3$$

is strictly increasing.

Proof. The assertions (a) and (b) (i)-(iii) were proved in [5], and (b) (iv) and (v) easily follow from the identity (see [1])

$$\frac{2}{\sqrt{\pi}} g(p,y) = \exp\left(\frac{p^2y^2}{4}\right)\left(\operatorname{erf}\left(\frac{py}{2}\right) + \operatorname{erf}\left(\frac{(2-p)y}{2}\right)\right).$$

Owing to function

$$\Phi_2(y) = \frac{y}{\sqrt{\pi}} \frac{\exp\left(-y^2\right)}{\operatorname{erfc}(y)} - y^2$$

is increasing (see [6],[7]) the assertion (c) holds. ■

In Section 2 we prove the existence and the uniqueness of the similarity solution for different values of the physical parameters ρ, β and d. In Section 3, we dicuss the equivalence of the problems I and II, and we extend some existence results for problem II obtained in [5].

2 Existence and Uniqueness of Similarity Solutions

In order to solve equation (10) we introduce the following function

$$Q_0(y) = \frac{\delta y + \nu y^3 + K_2 F(m,y)}{\exp\left((p-1)\,y^2\right)} \tag{13}$$

defined for $y > 0$ which verifies $Q_0(0) = K_2 A > 0$.

Theorem 3 *Let m be a positive real number. We define the following sets in the plane ν, p:*

$$R_1 = \left\{(\nu,p) \in \mathbb{R}^2 : -\sqrt{\pi}K_2\gamma_0 m \leq \nu,\ p \leq 1\right\}, \quad R_2 = \mathbb{R}^2 - R_1.$$

We have:

(a) If $(\nu,p) \in R_1$ then the problem I has a unique similarity solution if and only if $q_0 > \frac{K_F}{\alpha_2\sqrt{\pi}}A$.

(b) If $(\nu,p) \in R_2$ then the problem I has a similarity solution if and only if $0 < q_0 \leq \max_{y\geq 0} Q_0(y)$.

Proof. To prove the existence (and uniqueness) of similarity solution to problem I, it is necessary and sufficient to verify that the equation (10) has a (unique) solution. The equation (10) has a solution ξ if and only if $q_0 = Q_0(\xi)$. The proof is splitted in four cases [7]:

(i) $m > 0, \nu \geq 0$ and $p \leq 1$; (ii) $m > 0, \nu \geq 0$ and $p > 1$;

(iii) $\nu < -\sqrt{\pi}K_2\gamma_0 m$; (iv)$m > 0, -\sqrt{\pi}K_2\gamma_0 m \leq \nu < 0$ and $p \leq 1$. ■

Remark 4 *For $m = 0$, i.e. $d\rho = 0$, there exist a unique solution of equation (10) if and only if the inequality $q_0 > K_2 A$ is verified. This result has already been found in [10].*

Remark 5 *We note that in the case $p > 1$, if $\max_{y\geq 0} Q_0(y) > q_0 > K_2 A$ there exist at least two solutions. On the other hand, if q_0 is sufficiently small, then there exists a unique solution. The situation is a bit different in the problem II, studied in [5], where it was proved the existence and uniqueness of similarity solutions in the case $m > 0, \nu \geq 0, p \leq 2$.*

Similarly, we can obtain the following results.

Theorem 6 *Let $m < 0$. We define the sets*

$$R_3 = \left\{(\nu,p) \in \mathbb{R}^2 : \nu > -\sqrt{\pi}K_2\gamma_0 m,\ p \leq 1\right\}, \qquad R_4 = \mathbb{R}^2 - R_3.$$

Then

(a) If $(\nu,p) \in R_3$, there exists a solution when $q_0 > K_2 A$.

(b) If $(\nu,p) \in R_4$, there exists a solution when $0 < q_0 \leq \max_{y>0} Q_0(y)$.

Proof. By using Proposition 2 we have, $\delta y + \nu y^3 + K_2 F(m,y) \sim \nu + K_2\sqrt{\pi}\gamma_0 m y^3$ when $y \to +\infty$ and it follows that if $(\nu,p) \in R_3$ then $\lim_{y\to+\infty} Q_0(y) = +\infty$, from which we have $[K_2 A, +\infty) \subset \mathrm{Range}(Q_0)$. This proves part (a).

If $(\nu,p) \in R_4$, it is easy to see that the function Q_0 has a positive finite maximum, and then the second part is also proved. ■

Remark 7 *In [7] it was studied the physical acceptability of the similarity solution as a function of the thermal coefficients.*

3 Relationship between Problems I and II

Let (s, u, v) be given by (9), for some constant $\xi > 0$. Then $u(0, t)$ is a constant given by

$$u(0,t) = m\xi^2 + \frac{2q_0}{K_U} g(p,\xi) > 0. \tag{14}$$

Then, we can consider the problem II, by imposing this new temperature as $u(0,t)$ at the fixed face $x = 0$.

Theorem 8 *Let $m > 0$, $\nu > 0$, $p \leq 1$ and $q_0 > K_2 A$. If (s, u, v) is the unique similarity solution of Problem I, then (s, u, v) is the unique similarity solution of Problem II, provided the constant B in the condition (8) is given by*

$$B = m\xi^2 + \frac{2q_0\alpha_1}{K_U} g(p,\xi) \tag{15}$$

where ξ is the unique solution of equation (10).

Proof. We know that (s, u, v) is given by (9) where ξ is the unique solution of (10), which can be written as

$$Q_0(y) = q_0, \qquad y > 0 \tag{16}$$

By the results obtained in [5], there exists a unique solution to Problem II, with B defined by (15), given by

$$\begin{aligned} \bar{s}(t) &= 2\bar{\xi}\alpha_1\sqrt{t} \\ \bar{u}(x,t) &= B - \frac{m\bar{\xi} - B}{g(p,\bar{\xi})} \int_0^{\frac{x}{2\alpha_1\sqrt{t}}} \exp(p\bar{\xi}r - r^2)\, dr \\ \bar{v}(x,t) &= \frac{m\bar{\xi}^2 \operatorname{erfc}\left(\frac{x}{2\alpha_2\sqrt{t}}\right) + A\left(\operatorname{erf}\left(\gamma_0\bar{\xi}\right) - \operatorname{erf}\left(\frac{x}{2\alpha_2\sqrt{t}}\right)\right)}{\operatorname{erfc}\left(\gamma_0\bar{\xi}\right)} \end{aligned} \tag{17}$$

where $\bar{\xi}$ is the unique solution of the equation

$$\frac{\sqrt{\pi}}{2} K_1 \left(B - my^2\right) \frac{\exp\left((p-1)\,y^2\right)}{g(p,y)} - K_2 F(m,y) = \delta y + \nu y^3\,, \qquad y > 0 \tag{18}$$

and $K_1 = \frac{K_U}{\alpha_1 \pi}$. It is easy to see that the solutions given by (9) and (17) are coincident if and only if $\xi = \bar{\xi}$. Then, it is sufficient to see that ξ is a solution of (18) [7]. ∎

Suppose that (s, u, v) is a solution to problem I, with the boundary condition (7). By the results of Section 1, we know that (s, u, v) are given by (9), where ξ must satisfy the equation (10). For this solution, the temperature in the fixed boundary is constant and equal to $B = u(0,t) = T_0(q_0, \xi)$, where T_0 is the real function defined by

$$T_0(q,y) = my^2 + \frac{2}{\sqrt{\pi}} \frac{q}{K_1} g(p,y), \qquad q > 0,\, y > 0\,. \tag{19}$$

Assuming that $q > 0$, we will describe some properties of function T_0. First of all, we note that $T_0(q, 0) = 0$. Besides, it follows from the proposition 2 that if $m > 0$ and $p > 0$, then $T_0(q, y)$ is an increasing function in both of its arguments, with $T_0(q, +\infty) = +\infty$. If $m < 0$ and $p > 0$, then $T_0(q, +\infty) = +\infty$, and if $m < 0$ and $p \leq 0$, then $T_0(q, +\infty) = -\infty$. Finally, if $m > 0$ and $p < 0$ then $T_0(q, +\infty) = +\infty$.

Suppose that $m > 0$. For each $\bar{\xi} > 0$ let

$$\bar{q}_0 = Q_0(\bar{\xi}) \tag{20}$$

where Q_0 is the function defined by (13). Let $\bar{B} = T_0(\bar{q}_0, \bar{\xi}) = m\bar{\xi}^2 + \frac{2}{\sqrt{\pi}} \frac{\bar{q}_0}{K_1} g(p, \bar{\xi})$, then a solution to problem I with $q_0 = \bar{q}_0$, which is given by (9) with $\xi = \bar{\xi}$ because of (20), corresponds to a solution to Problem II with $B = \bar{B}$. Then, given $B > 0$, we can show the existence of solution to Problem II, by proving that B belongs to the image set of the function $J(\cdot) = T_0\left(Q_0\left(\cdot\right), \cdot\right)$. For different values of the physical parameters, we study function J and we obtain the following results [7].

Theorem 9 *Let $m > 0$. If $\nu \geq -\sqrt{\pi} K_2 m \gamma_0$, then there exists a similarity solution to Problem II. If, in addition, $0 \leq p \leq 1$, then the similarity solution is unique. For the case $\nu < -\sqrt{\pi} K_2 m \gamma_0$, a sufficient condition in order to have the existence of a solution to Problem II, is that B verifies the inequality*

$$B < m\frac{\delta}{|\nu|} + \frac{2}{\sqrt{\pi}}\frac{M}{K_1} g\left(p, \sqrt{\frac{\delta}{|\nu|}}\right).$$

Remark 10 *The last Theorem, extends a result of [5], where it was proved that if $m > 0, \nu < 0$, then there exists a solution to Problem II when*

$$B < \frac{m\delta}{|\nu|}.$$

Proposition 11 *Suppose $m < 0$ and $\nu > \sqrt{\pi} K_2 |m| \gamma_0$, then:*

(a) If $p \leq 2$ the problem II has a similarity solution if

$$B > \max(0, J(y_0)) \tag{21}$$

(b) If $p > 2$ the problem II has a similarity solution if $m + \frac{2(\nu + \sqrt{pi} K_2 m \gamma_0)}{\sqrt{pi} K_1 (p-2)} > 0$ and (21) are verified.

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